

EMERGENCE OF SELF-SUSTAINED OSCILLATIONS  
IN A HYPOTHETICAL METABOLIC PATHWAY REGULATED  
BY A POSITIVE FEEDBACK LOOP

*Aparición de oscilaciones autosostenidas en una ruta metabólica  
hipotética regulada por un bucle de retroalimentación positiva*

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Recibido: 14/10/2020 • Aprobado: 28/10/2020

**Cómo citar:** Boluda, C. J., Ceballos, M., Díaz, K., Hernández, D., Hidalgo, D., & Martínez, P. (2020). Aparición de oscilaciones autosostenidas en una ruta metabólica hipotética regulada por un bucle de retroalimentación positiva. *Ciencia, Ingenierías y Aplicaciones*, 3(2), 5-38. Doi: <https://doi.org/10.22206/cyap.2020.v3i2.pp5-38>

**Abstract**

*Sustained oscillatory behaviors can emerge in both chemical and biological systems, although more frequently in the latter. In metabolic pathways, periodic oscillations to which functional importance is attributed to have been observed with periods of the order of minutes. This biochemical rhythmicity can be explained through the existence of an allosteric enzyme, which is considered the origin of the system's structural instability, promoting its evolution towards non-stationary dynamics such as sustained oscillations and chaos. This article presents a simple kinetic model of a hypothetical 4-step linear metabolic pathway with a positive feedback loop. The model allows the simulation of periodic behaviors for analyzing their triggering factors and underlying mechanisms. The oscillations arise due to the positive feedback for a narrow range of the inlet flux. The amplitude of the oscillations in metabolite concentrations depends on the oscillator's distance, kinetic constants' value, and the inlet flux. Under forced oscillatory inlet fluxes, more complex behaviors arise.*

**Keywords:** metabolic pathway; non-stationary dynamic; positive feedback loop; sustained oscillations.

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## Resumen

*Los comportamientos oscilatorios sostenidos pueden surgir tanto en sistemas químicos como biológicos, aunque con mayor frecuencia en estos últimos. En las vías metabólicas se han observado oscilaciones con periodos del orden de minutos a las que se atribuye importancia funcional. Esta ritmicidad bioquímica puede explicarse mediante la existencia de una enzima alostérica que es considerada el origen de la inestabilidad estructural del sistema, promoviendo su evolución hacia dinámicas no estacionarias como las oscilaciones sostenidas y el caos. Este artículo presenta un modelo cinético simple de una vía metabólica hipotética lineal de cuatro etapas y que consta de un bucle de retroalimentación positiva. El modelo permite la simulación de comportamientos periódicos con el objeto de poder analizar sus factores desencadenantes y mecanismos subyacentes. Las oscilaciones surgen debido a la retroalimentación positiva para un rango estrecho de valores del flujo de entrada. La amplitud de las oscilaciones en las concentraciones de metabolitos depende de la distancia al oscilador, el valor de las constantes cinéticas y el flujo de entrada. El sistema exhibe comportamientos más complejos cuando el flujo de entrada es oscilatorio.*

**Palabras clave:** dinámicas no estacionarias; retroalimentación positiva; ruta metabólica; oscilaciones sostenidas.

## 1. Introduction

### 1.1 Oscillations in chemical and biological systems

In the early 1950s, Russian biochemist Boris Pavlovich Belousov discovered that the oxidation of citric acid, with bromate ( $\text{BrO}_3^-$ ) and cerium ions as a catalyst, did not progress as expected. Instead of a gradual change from a colorless solution to a pale yellow one, due to the oxidation of  $\text{Ce}^{3+}$  to  $\text{Ce}^{4+}$ , periodic fluctuations in the color of the reaction mixture were observed, indicating that successive interconversions between reduced and oxidized states of reactants were taking place. Belousov was unsuccessful in publishing his results because the reviewers believed that this phenomenon disagreed with the Second Law of Thermodynamics (Kiprijanov, 2016). However, in 1961 Zhabotinsky verified Belousov's discovery, publishing his results in 1964 (Tyson, 1994). The Belousov-Zhabotinsky reaction is considered a prototype of an oscillatory chemical reaction, being its mechanism well-known (Noyes et al., 1972). Nevertheless, oscillations in chemical systems were discovered long before. As early as the 17<sup>th</sup> century, Robert Boyle observed the intermittent ignition of phosphorus (Agreda & Barragán, 1998), and references from the 19<sup>th</sup> century mention the mercury beating heart, an electrochemical-based mechanical oscillator (Martín & Martín, 2010). In 1910 Lotka proposed a hypothetical reaction model that displays damped oscillations (Lotka, 1910).

Later on, in 1920, the same author showed that chemical systems could adopt sustained oscillatory behavior under certain circumstances (Lotka, 1920). However, it was not until the late 1960s when other researchers became interested in these types of reactions (Goldbeter, 2018). In 1956 Prigogine and Balescu demonstrated that sustained oscillations could occur in open chemical systems with nonlinear kinetics and far from equilibrium (Prigogine & Balescu, 1956; Goldbeter, 1996). In the 1960s, Goodwin studied negative feedback as a source of oscillations and the quasiperiodic and subharmonic behavior resulting from the interaction between two oscillators. Goodwin ventured the importance of these phenomena for cellular physiology (Gonze & Ruoff, 2020). In 1967 Prigogine and Nicolis demonstrated that sustained oscillations do

not contradict any thermodynamic principle and showed that the instability of these chemical oscillators is due to positive (PFBLs) or negative feedback loops (NFBLs) (Prigogine & Nicolis, 1967; Agreda & Barragán, 1998).

In 1968 Lefever, Nicolis and Prigogine presented an autocatalytic theoretical model known as Brusselator (an acronym for Brussels and oscillator) that has been very useful for the understanding of oscillatory behaviors in chemical systems (Lefever et al., 1967; Goldbeter, 1996; Agreda & Barragán, 1998). In contrast to the scarcity of oscillatory behavior in chemical systems, rhythmicity is inherent to life (Li & Yang, 2018). Although periodic phenomena are not exclusive to biological systems, this behavior is much more frequent in living beings than in purely chemical systems (Goldbeter, 1996; Goldbeter, 2018). The circadian rhythm of sleep and wakefulness is a classic example of rhythmicity in which both frequency and amplitude of a circadian oscillator must keep robust 24-hour rhythms in the presence of external fluctuations (Kang & Cho, 2017). Although oscillations are necessary for many biological functions that require precise control of time, biological systems are inherently noisy, affecting its oscillations' accuracy.

However, these systems achieve robustness by dissipating a critical amount of free energy, which results in greater irreversibility and accuracy of the oscillations (Cao et al., 2015; Fei et al., 2018). The relationship between the amount of energy dissipated and the precision of the oscillations period is a particularly topical question that has been the subject of recent research (del Junco and Vaikuntanathan, 2020). Heart and respiratory rates, fluctuations in hormone levels, electrically excitable cells like neurons, or muscle cells are additional examples of rhythmicity. Physiological oscillations can be a consequence of the central nervous system's activity, although in certain instances, they are generated *in situ*. Thus, the heart rate is due to the sinoatrial node's pacemaker cells, located in the right atrium, which have an unstable resting membrane potential.

Rhythmically and spontaneously, these cells generate an action potential that is transmitted with the necessary delay to coordinate the atria and ventricles' function throughout the myocardium whose cells are

electrically coupled. The frequency of the oscillations can become high; for example, in those insects equipped with asynchronous flying muscles, as *Diptera*, the wings can reach, in certain species, a frequency of 1000 Hz (Ruppert & Barnes, 1996). Oscillatory behavior is also observed in non-excitable cells, such as intracellular  $\text{Ca}^{2+}$  fluctuations, enzymatic reactions, or the secretion of many hormones released from the endocrine glands in a pulsatile manner rather than continuously (Goldbeter, 2018). Additional examples of oscillatory behaviors in non-excitable cells are the nuclear factor kappa B (NF- $\kappa$ B), and the proteins p53 and p38 (Li & Yang, 2018). The latter is a mitogen-activated protein kinase (MAPK), which controls an inflammatory response; oscillations in its activity have been observed for more than 8 hours under stimulation by interleukin-1 $\beta$  (Tomida et al., 2015).

The tumor suppressor protein p53, regulated through a negative feedback loop (NFBL) with the *mdm2* protein, also exhibits an oscillatory dynamic. This mechanism allows the cell to repair its DNA through the action of p53 without excessive exposure to said protein, which could induce the apoptotic cell-death (Bar-Or et al., 2000). The oscillations of p53 are observed in cell lines of different species, although with different frequencies. For example, in rodent cells, these oscillations are faster than those observed in cell lines derived from humans, monkeys, and dogs (Stewart-Ornstein et al., 2018). Nuclear factor-kappa B (NF- $\kappa$ B), a protein complex that plays a crucial role in inflammation and cancer, exhibits damped oscillations synchronized with external disturbances (Zambrano et al., 2016).

## 1.2 Biochemical oscillations

The study of oscillating biochemical systems began in 1955 after observing changes in the concentrations of 3-phosphoglycerate and ribulose-1,5-diphosphate during the dark phase of photosynthesis (Chance et al., 1973). Oscillations in glycolysis were observed in 1957, while continuous glycolytic oscillations in cell-free yeast systems were demonstrated in 1966. (Chance et al., 1973; Agreda & Barragán, 1998; Robles & Barragán, 2006). Oscillations of cAMP (Li & Goldbeter, 1990; Goldbeter, 2018), intracellular  $\text{Ca}^{2+}$  (Dupont & Goldbeter, 1992), and glycolysis metabolites (Goldbeter, 1996) are among the best-known

examples of periodic behavior at the biochemical level. The role of oscillations can be diverse and have high functional importance; abnormalities in them can cause disease, including cancer (Li & Yang, 2018). It has been suggested that they may be at the origin of some circadian rhythms and other processes such as cell division (Rodríguez et al., 1998). Thus, in *Dyctiostelium discoideum* amoebas, concentration's periodic oscillations are related to the aggregation of cells that respond chemotactically to cAMP pulses (Goldbeter 1996; Goldbeter, 2018).

Periodic and chaotic behaviors are observed in metabolic pathways driven by matter flows and controlled by positive or negative feedback loops (PFBLs or NFBLs) (Goldbeter, 1996). In the context of biochemical pathways, feedback loops (FBLs) consist of at least two metabolites that regulate their activity reciprocally by activation or inhibition. The complex behavior of biological systems results from the coexistence of multiple of these FBLs (Dalchau et al., 2018).

Glycolysis is considered the central route of glucose catabolism and provides cells a significant portion of energy and pyruvate that can continue the subsequent oxidative degradation. It is a 10-steps metabolic pathway that transforms glucose and other hexoses into two pyruvate molecules, obtaining two molecules of ATP and NADH. This metabolic pathway begins and ends with irreversible steps: the conversion of glucose to glucose-6-phosphate by hexokinase, consuming ATP (step 1), and the conversion of phosphoenolpyruvate to pyruvate by pyruvate kinase, producing ATP (step 10). There is a third irreversible step: the transformation of fructose-6-phosphate to fructose-1,6-bisphosphate by phosphofructokinase (PFK). These three irreversible steps suppose a greater dissipation of energy, driving glycolysis away from thermodynamic equilibrium (Robles & Barragán, 2006).

Metabolic oscillations in this pathway have been observed in tumor cells, muscle extracts, fibroblasts, and yeast, the latter being the most studied system (Sel'kov, 1968; Robles & Barragán, 2006). These oscillations emerge in glycolysis when glucose-6-phosphate or fructose-6-phosphate is supplied, but not for the addition of fructose-1,6-diphosphate. This observation suggests that the enzyme phosphofructokinase (PFK) is the oscillophore responsible for the periodic behavior of the pathway,

which can also be demonstrated by adding inhibitors or activators of the enzyme. Specifically, the addition of ADP (an activator of PFK) affected the frequency and amplitude of the oscillations of glycolytic metabolites; this is an example of PFBL.

Although PFBLs exerted by a product of a metabolic pathway on an enzyme are not as frequent as product inhibition, autocatalysis has an essential role in biochemical oscillations (Robles & Barragán, 2006; Tyson, 2002). Nonetheless, the importance of PFK in the oscillations observed in this metabolic pathway has also been relativized, proposing that the characteristics of the oscillations also depend on other enzymes of the pathway (Reijenga et al., 2005). Alternatives to the oscillophore have been proposed to explain this glycolytic behavior through the autocatalytic character of the stoichiometry of the pathway, which acts as a PFBL responsible for the oscillations (Gustavsson et al., 2014).

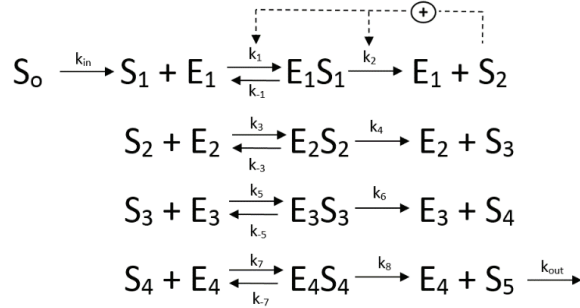
The theoretical approach to these periodic biochemical phenomena through models is useful in understanding their physiological roles and underlying mechanisms. An example of the latter is the Oregonator, an autocatalytic model that reproduces the oscillatory behavior of the Belousov-Zabotinsky reaction (Field & Noyes, 1974).

This article presents a hypothetical metabolic pathway with a PFBL. The system is kept away from thermodynamic equilibrium by constant input and output fluxes. The positive feedback consists of the activation of the first enzyme of the pathway by its product. Our main target is to provide a simple kinetic model in which, intrinsically and under certain conditions, sustained oscillatory behavior arises. Such a model could be a starting point for building modified or more complex and realistic models for educational and research purposes.

## 2. Model and kinetic equations

Figure 1 shows the sequence of enzymatic reactions of a hypothetical linear metabolic pathway consisting of four steps and six metabolites:  $S_0$ ,  $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ . The concentration of  $S_0$  is constant and it is also converted at a constant rate by a non-enzymatic irreversible step in  $S_1$ , representing a non-rhythmic input of the substrate into the system. The

rest of the metabolites are sequentially transformed from  $S_1$  to  $S_5$  through the subsequent action of the enzymes  $E_1$ ,  $E_2$ ,  $E_3$ , and  $E_4$ .



**Figure 1.** Sequential reactions that make up a linear metabolic pathway

**Note:** The conversion from  $S_1$  to  $S_2$  consists of a positive feedback loop (PFBL), thus being a self-catalytic step. Dashed arrows represent the activation of the step by  $S_2$ .

As described in Figure 2, each enzyme reversibly binds to its substrate in these four enzymatic reactions, forming a substrate-enzyme complex ( $E_x S_x$ ) that converts into the corresponding product. The kinetic constants for the reversible formation of the substrate enzyme complex are given by  $k_n$  and  $k_{-n}$ . For each step of the pathway, the substrate enzyme complex ( $E_x S_x$ ) yields its corresponding product ( $S_{x+1}$ ) with a kinetic constant denoted by  $k_{n+1}$  and  $k_{(n+1)-1}$  for the reverse reaction. When  $k_{(n+1)-1} = 0$ , the conversion of  $E_x S_x$  into  $E_x$  and  $S_{x+1}$  is irreversible.



**Figure 2.** Enzyme reversibly binds its substrate forming a substrate-enzyme complex

**Note:** Each enzyme of the pathway reversibly binds its substrate to form an enzyme-substrate complex ( $E_x S_x$ ) with  $k_n$  and  $k_{-n}$  as kinetic constants of the direct and reverse reaction. The  $E_x S_x$  complexes evolve to form their corresponding products ( $S_{x+1}$ ) with a constant rate of  $k_{n+1}$ .

The dynamic of the enzymatic conversions showed in Figure 1 can be described by the following differential equations:



$$\begin{aligned}
1) \quad \frac{d[S_1]}{dt} &= k_{in}[S_0] - k_1(1 + [S_2]^3)[E_1][S_1] + k_{-1}[E_1S_1] \\
2) \quad \frac{d[S_2]}{dt} &= k_2(1 + [S_2]^3)[E_1S_1] - k_3[E_2][S_2] + k_{-3}[E_2S_2] \\
3) \quad \frac{d[S_3]}{dt} &= k_4[E_2S_2] - k_5[E_3][S_3] + k_{-5}[E_3S_3] \\
4) \quad \frac{d[S_4]}{dt} &= k_6[E_3S_3] - k_7[E_4][S_4] + k_{-7}[E_4S_4] \\
5) \quad \frac{d[S_5]}{dt} &= k_8[E_4S_4] - k_{out}[S_5]
\end{aligned}$$

As shown in Figure 1, the  $E_1$ -catalyzed transformation of  $S_1$  to  $S_2$  simulates an activation per product step. The formation reaction rate for the  $E_1S_1$  complex and its conversion into  $E_1$  and  $S_2$  have become strongly dependent on the concentration of  $S_2$  by the inclusion of the algebraic expression  $1 + [S_2]^3$  in equations 1 and 2. This represents an activation of the  $E_1$  enzyme by the substrate  $S_2$ , increasing both the affinity of the enzyme  $E_1$  to the substrate  $S_1$  and the conversion rate of this complex to  $S_2$  and  $E_1$ . The inclusion of the binomial  $1 + [S_2]^3$  in the differential equations is responsible for the system's non-linearity. According to the Arrhenius equation, the kinetic constants depend not only on the temperature but also on the presence of catalyst; therefore, it can be considered that the expression  $(1 + S_2^3)$  modifies the kinetic constants as follows:  $k_1' = k_1(1 + S_2^3)$  and  $k_2' = k_2(1 + S_2^3)$ .

Since in living cells, the inlet flow is considered a key parameter to trigger the oscillations (Danø et al., 1999); the model's behavior was evaluated for concentration values of  $S_0$  between 0 and 12 arbitrary units (a.u.). In all experiments, the inlet flux was such that the concentration of  $S_2$  (activator of  $E_1$ ) was lower than the concentration of  $E_1$  ( $[S_2] < [E_1]$ ). Table 1 shows the values of the kinetic parameters, concentration of enzymes, and initial substrates concentration for each of the pathway steps.

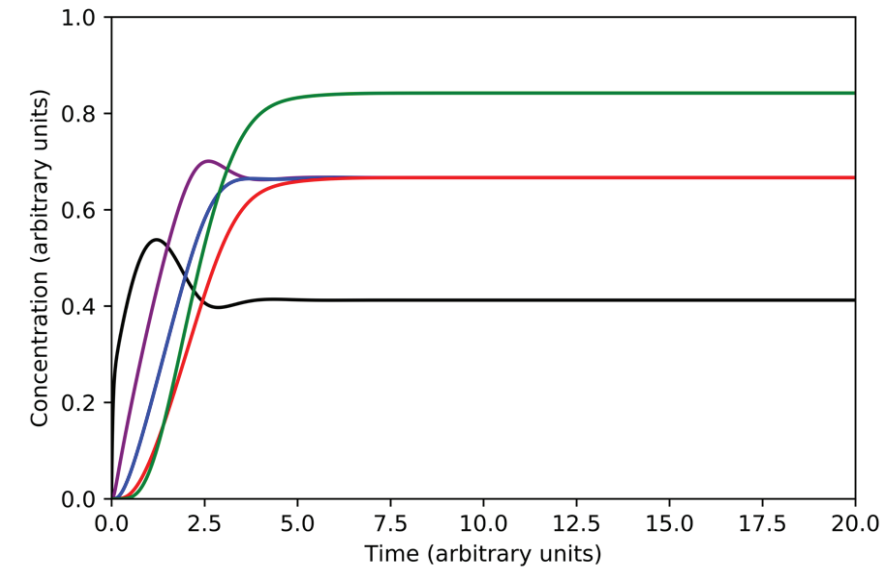
**Table 1.** Metabolites and enzyme’s concentration (a.u.) and kinetic parameters

Metabolites concentration (arbitrary unit)	Enzymes Concentration (arbitrary unit)	Kinetic constants
$S_0=0-12$	-	$k_{in}=1$
$S_1=0$	$E_1=10$	$k_1=3$
		$k_{-1}=1$
		$k_2=2$
		$k_{-2}=0$
$S_2=0$	$E_2=10$	$k_3=3$
		$k_{-3}=1$
		$k_4=2$
		$k_{-4}=0$
$S_3=0$	$E_3=10$	$k_5=3$
		$k_{-5}=1$
		$k_6=2$
		$k_{-6}=0$
$S_4=0$	$E_4=10$	$k_7=3$
		$k_{-7}=1$
		$k_8=2$
		$k_{-8}=0$
$S_5=0$	-	$k_{out}=9.5$

Python version 3.7.4 with Spyder and Anaconda as IDE was used for numerical integration of differential equations. The simulations were performed, implementing the kinetic parameters and concentration values for metabolites and enzymes shown in Table 1. Python script used for the experiments is provided in the supplementary materials.

3. Results

Using the values specified in Table 1 and for [So] equal to 8 a.u., the system reaches the steady state (Figure 3) with saturation percentages for enzymes not exceeding 40%. The steady state concentrations of metabolites and enzyme saturation percentages are shown in Table 2.



**Figure 3.** Evolution of metabolite concentrations ( $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ ) as the system reaches the steady state when  $[S_0] = 8$  a.u.

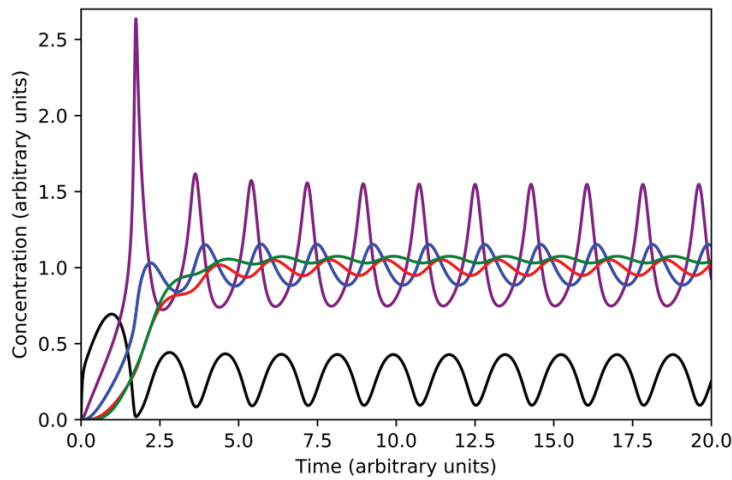
**Table 2.** Metabolite concentrations ( $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ ) and enzyme saturation percentages when the system reaches steady state

Metabolite	Concentrations in steady state (arbitrary unit)	Enzyme saturation percentage
$S_1$	0.4123	$E_1 = 30.86$
$S_2$	0.6667	$E_2 = 40$
$S_3$	0.6667	$E_3 = 40$
$S_4$	0.6667	$E_4 = 40$
$S_5$	0.8421	-

**Note:** In this experiment  $[S_0] = 8$  a.u.

However, when  $[S_0] = 10$  a.u., the system does not reach a steady state but is installed in a sustained oscillation, as displayed in Figure 4. Periodic oscillations with  $T=1.775$  a.u. are observed in all metabolites

of the pathway. The transformation of  $S_1$  to  $S_2$  speeds up as  $S_2$  increases until  $S_1$  depletion (Figure 1), limiting  $S_2$  production to a value that allows recovery of  $S_1$  levels, this being a classic oscillation generator mechanism (Tyson, 2002).



**Figure 4.** Undamped oscillations of the metabolites ( $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ ) concentrations are observed when  $[S_0]$  equals 10 a.u.

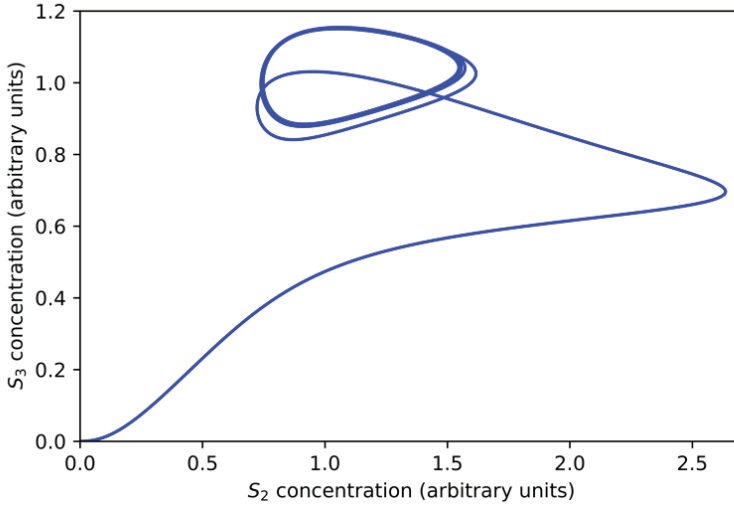
The oscillatory ranges of the concentration of metabolites are shown in Table 3.

**Table 3.** Oscillation intervals and amplitudes for metabolites when  $S_0=10$  a.u.

Metabolite	Oscillation amplitude	Oscillation interval
$S_1$	0.33287	0.09472-0.42759
$S_2$	0.79978	0.74726-1.54704
$S_3$	0.26731	0.88438-1.15169
$S_4$	0.10407	0.94837-1.05244
$S_5$	0.04433	1.03008-1.07441

**Note:** The amplitude of the oscillation is maximum for the metabolite  $S_2$ , which acts as the activator of  $E_1$  and minimum for  $S_5$ , which is the substrate that is further from the oscillations source.

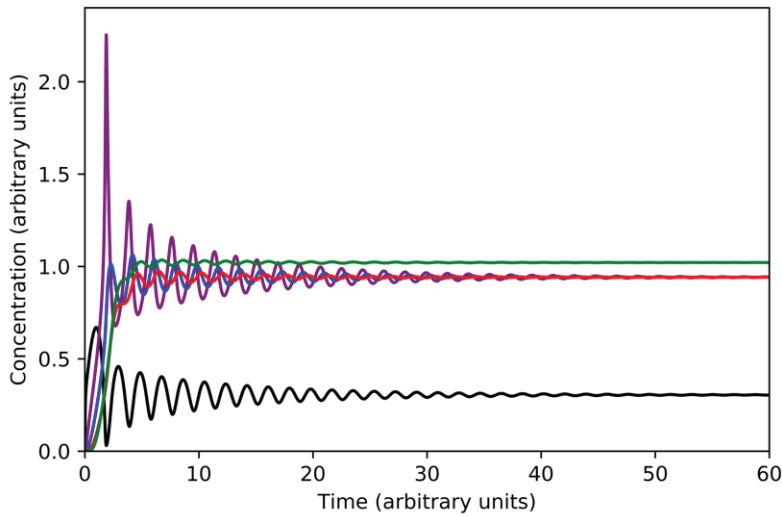
Figure 5 illustrates the phase plane plot for  $S_2$  against  $S_3$ , observing a limit cycle that is the proper attractor for the oscillatory behavior.



**Figure 5.** Phase plane plot of  $S_3$  vs.  $S_2$  when  $[S_0] = 10$  a.u., a limit cycle is observed

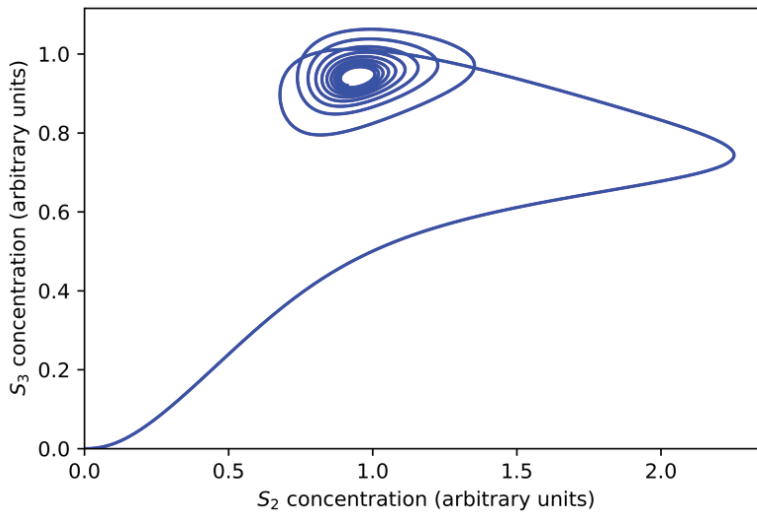
The transition from the steady state to the sustained oscillation takes place through a Hopf bifurcation (Gustavsson et al., 2014; Beta & Kruse, 2017). For a critical value of  $S_0$  (approximately 9.8 a.u.), the steady state becomes unstable, and the system evolves to an oscillatory behavior. When  $S_0 = 9.7$  a.u., near the bifurcation point, the system shows damped oscillations, which eventually reaches the steady state as shown in Figure 6. Figure 7 shows the phase plane for this experiment; a stable focus is attained. Figure 8 displays a close-up view of the said attractor.

(B)



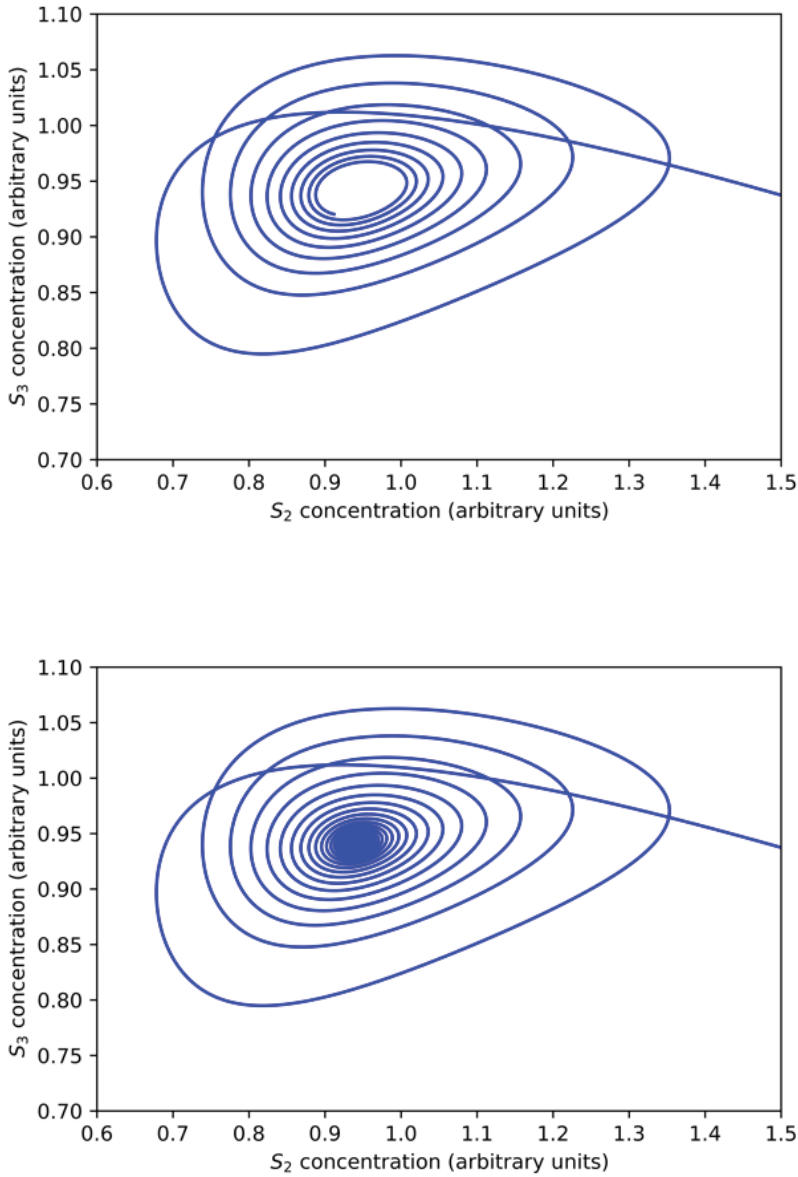
**Figure 6.** Damped oscillations of metabolites ( $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ ) around a fixed point for a value of  $[S_0] = 9.7$  a.u.

**Note:** The system slowly reaches steady state as the oscillations attenuate.



**Figure 7.** Phase plane plot of  $S_2$  vs.  $S_3$  when the concentration of  $S_0$  is equal to 9.7 a.u.

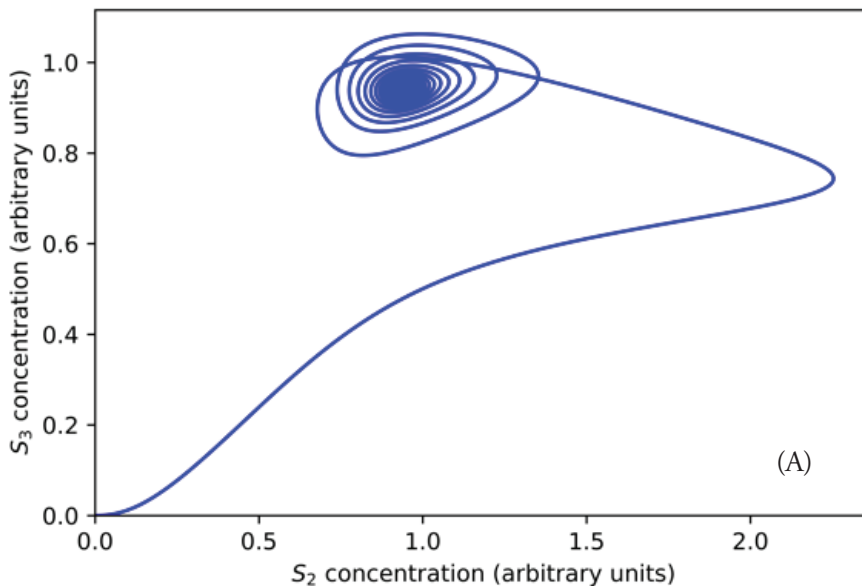
**Note:** The system evolves towards a fixed-point attractor.



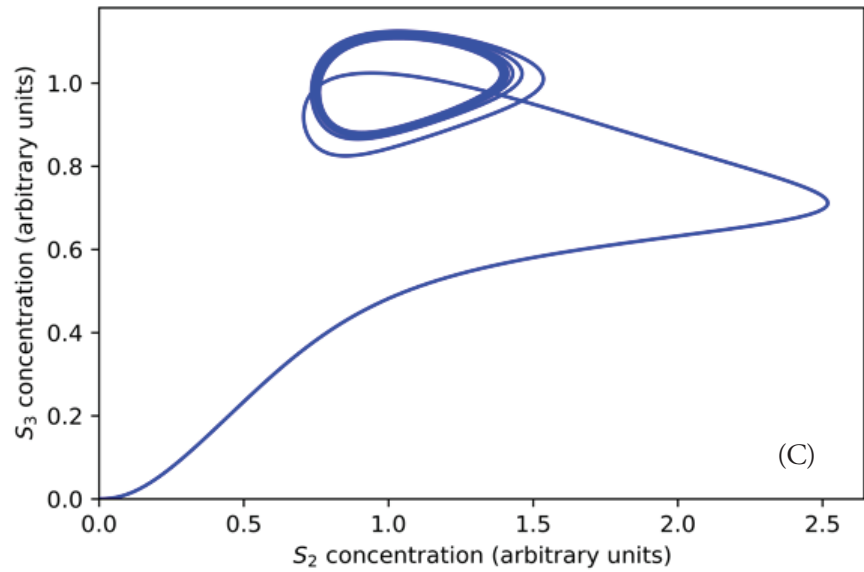
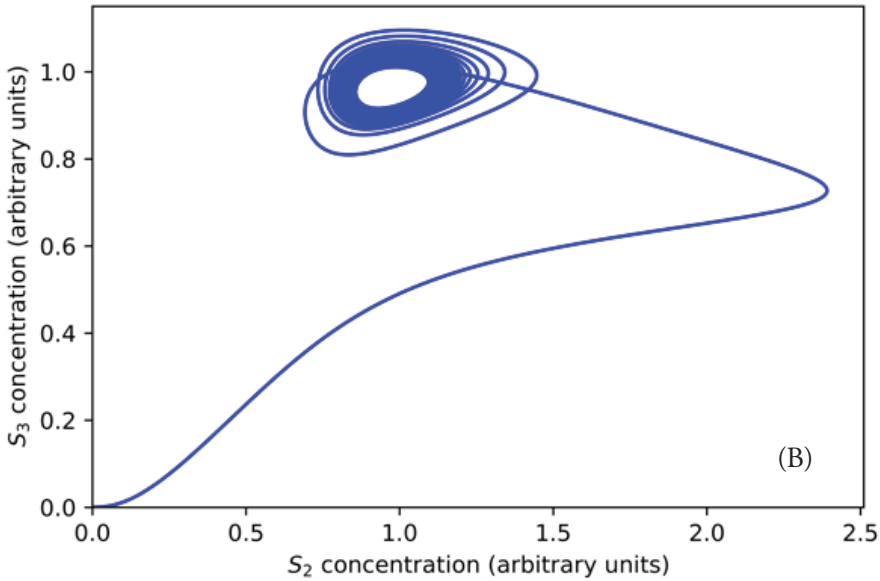
**Figure 8.** Close-up view of Figure 7

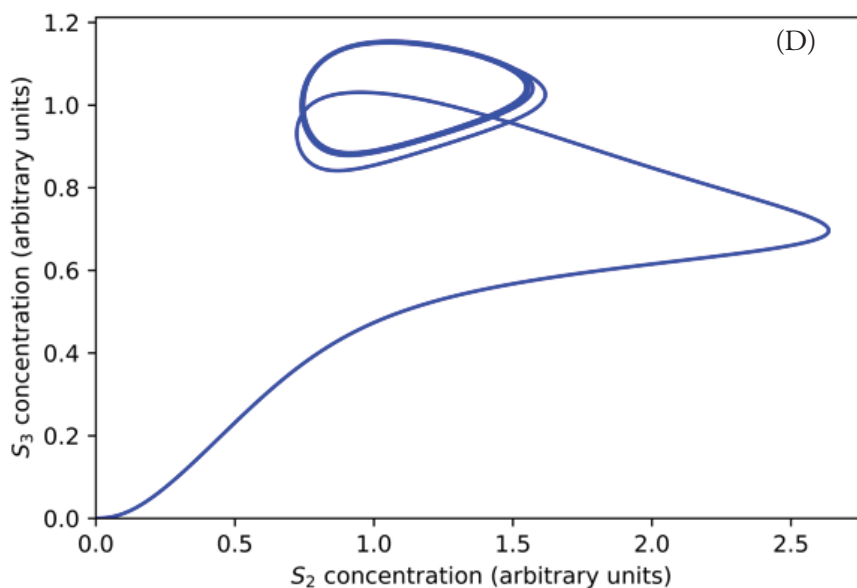
**Note:** When the concentration of  $S_0$  is equal to 9.7 a.u., the system evolves to a fixed-point attractor through damped oscillations. As time increases, the system gets closer to the steady state, as shown in the figure on the right.

The system's evolution from a fixed-point attractor towards the limit cycle has also been studied (Figure 9). As previously explained, when the concentration of  $S_0$  is equal to 9.7 a.u. the limit cycle is not achieved, resulting in a stable focus (Figure 9A). However, starting from a concentration of  $S_0$  equal to 9.8 a.u., a stable limit cycle starts taking shape as a result of the self-sustained oscillations, as described by Saha & Gangopadhyay (Saha & Gangopadhyay, 2017). This tendency is confirmed for concentration values of  $S_0$  equal to 9.9 and 10 a.u., for which their corresponding phase planes are presented in Figure 9, labeled as (C) and (D). The system returns to a steady state when the concentration of  $S_0$  is equal or higher than 11.75 (a.u.), as demonstrated in Figure 10. It should be clarified that the emergence of oscillatory behavior also depends on the concentrations of the enzymes. For example, changing the concentrations of  $E_1$ ,  $E_2$ ,  $E_3$ , and  $E_4$  to 15 a.u. and  $[S_0]$  equal to 10 a.u., the system does not oscillate. Instead, it reaches a steady state, thus making the oscillatory window even narrower. Contrary to what happens within the oscillatory interval, the concentration of  $S_2$  did not drop enough to allow a significant increase in  $S_1$ , which remains at low concentration, being this the probable reason for the lack of oscillations.



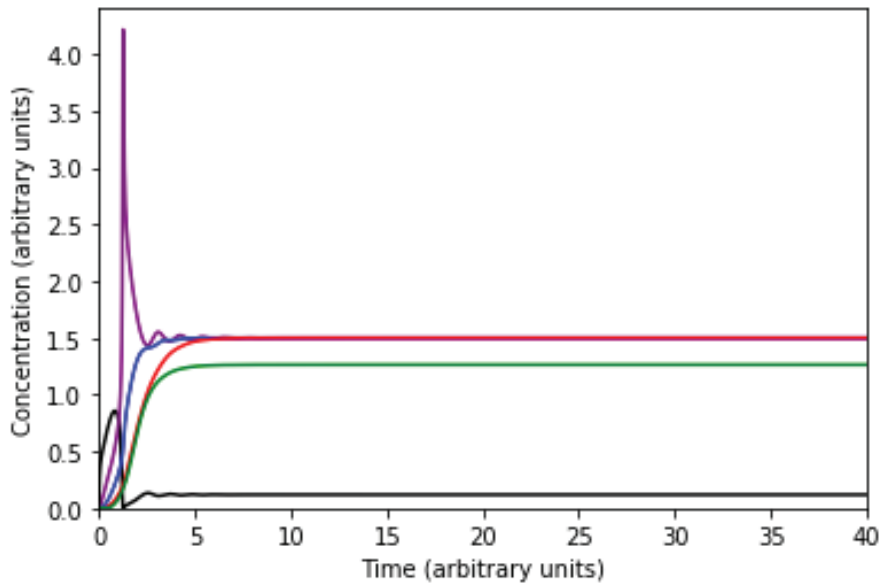






**Figure 9.** Four phase planes plot for  $S_2$  vs.  $S_3$  displays a fixed-point attractor's evolution towards a limit cycle.

**Note:** (A): When  $[S_0]$  is equal to 9.7 a.u., a stable focus is observed. (B) and (C) show the progressive trend to a limit cycle when  $[S_0]$  is equal to 9.8 a.u. and 9.9 a.u., respectively. When  $[S_0]$  is equal to 10 a.u., a defined limit cycle is observed (D).

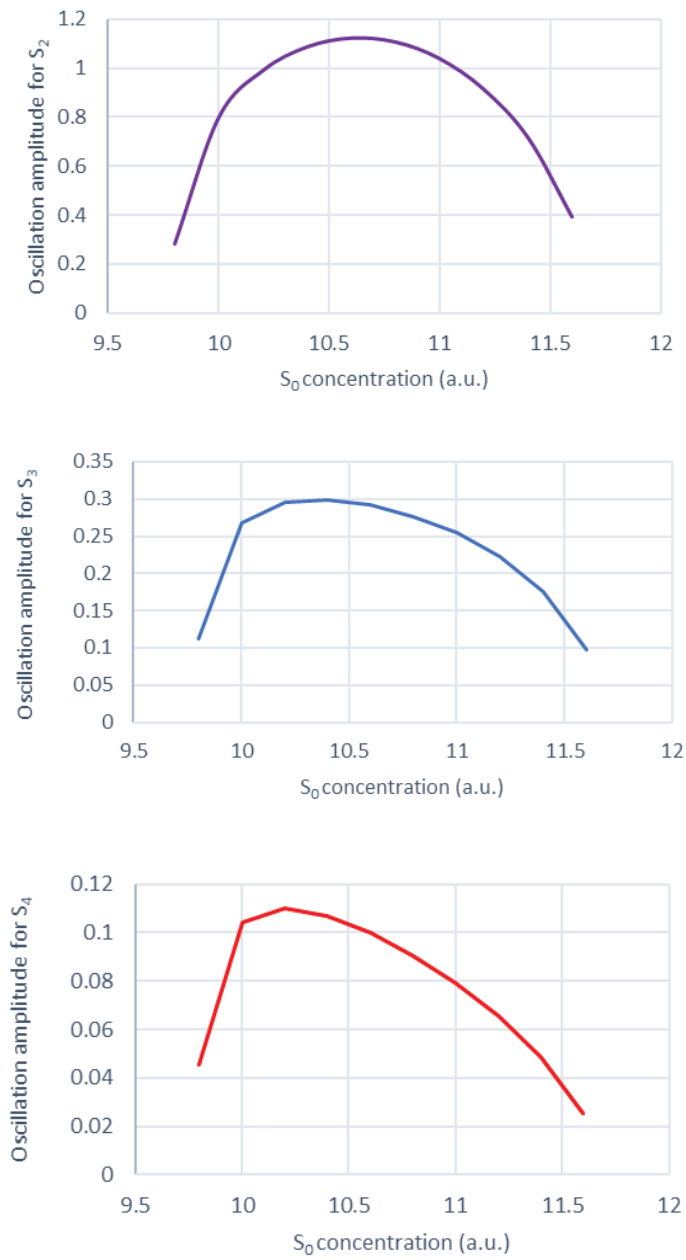


**Figure 10.** Evolution of metabolite concentrations ( $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ ) as the system reaches steady state for  $[S_0]=12$  a.u.

**Note:** The concentration of metabolites becomes constant over time.

Table 4 shows the amplitude of metabolites' oscillation as a function of  $[S_0]$ . Different metabolites reach maximum oscillations for different values of  $S_0$ . As displayed in Figure 11, oscillations' amplitude decreases progressively until the system returns to a steady state. Sustained oscillatory behavior is observed in a narrow interval of  $S_0$  values (approximately 9.8-11.6 a.u.). Outside this interval, the system remains in steady state, which is reached through damped oscillations near the oscillatory behavior's emergence point.

**Table 4.** Amplitude of metabolite oscillations for  $[S_0]$  values between 9.8 and 11.60 a.u.



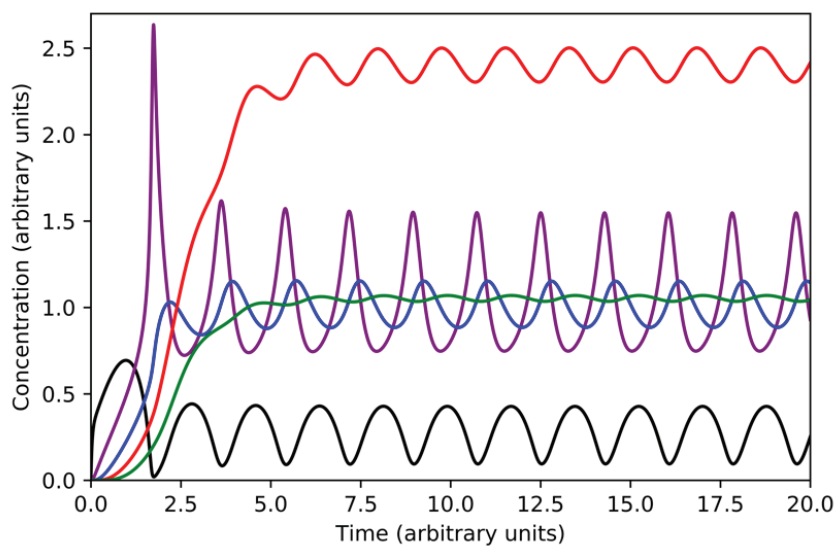
**Note:** For higher or lower  $[S_0]$  values, the system reaches a steady state.

[S <sub>0</sub> ] (a.u.)	Oscillation amplitude (a.u.)				
	[S <sub>1</sub> ]	[S <sub>2</sub> ]	[S <sub>3</sub> ]	[S <sub>4</sub> ]	[S <sub>5</sub> ]
9.8	0.15069	0.27998	0.11205	0.04551	0.02028
10	0.33287	0.79978	0.26731	0.10407	0.04433
10.2	0.35877	0.99127	0.29489	0.10989	0.04434
10.4	0.35818	1.08793	0.29857	0.10662	0.04068
10.6	0.34496	1.12382	0.29154	0.09969	0.03596
10.8	0.32221	1.10776	0.27680	0.09043	0.03086
11	0.29010	1.03964	0.25424	0.07909	0.02553
11.2	0.24721	0.91381	0.22192	0.06542	0.01997
11.4	0.18982	0.71660	0.17522	0.04866	0.01402
11.6	0.10242	0.39119	0.09752	0.02534	0.00688

**Figure 11.** Evolution of the amplitude of the metabolite’s oscillations (S<sub>2</sub>, S<sub>3</sub>, S<sub>4</sub>, and S<sub>5</sub>) as a function of S<sub>0</sub> concentration

The effect of the reversibility of the pathway’s steps was also briefly analyzed. Setting the value of  $K_{-8}$  to 2, instead of the previously stated value of  $k_{-8}=0$ , the conversion of  $E_4S_4$  into  $S_5$  becomes reversible (Figure 2). This means  $E_4$  can bind  $S_5$  to regenerate the  $E_4S_4$  complex that could in turn dissociate, leading to a higher concentration of  $S_4$ . As can be seen in Figure 12, when the last step of the pathway is reversible, the oscillations amplitude of  $S_4$  have a wider interval than those shown in Figure 4, which corresponds to a simulation where the last step is irreversible. The concentration intervals and its oscillation amplitudes of these metabolites are shown in Table 5.

This experiment demonstrates that the oscillation’s characteristics will depend not only on the oscilophore, but also on the kinetic parameters of other metabolic pathway steps.



**Figure 12.** Sustained oscillations

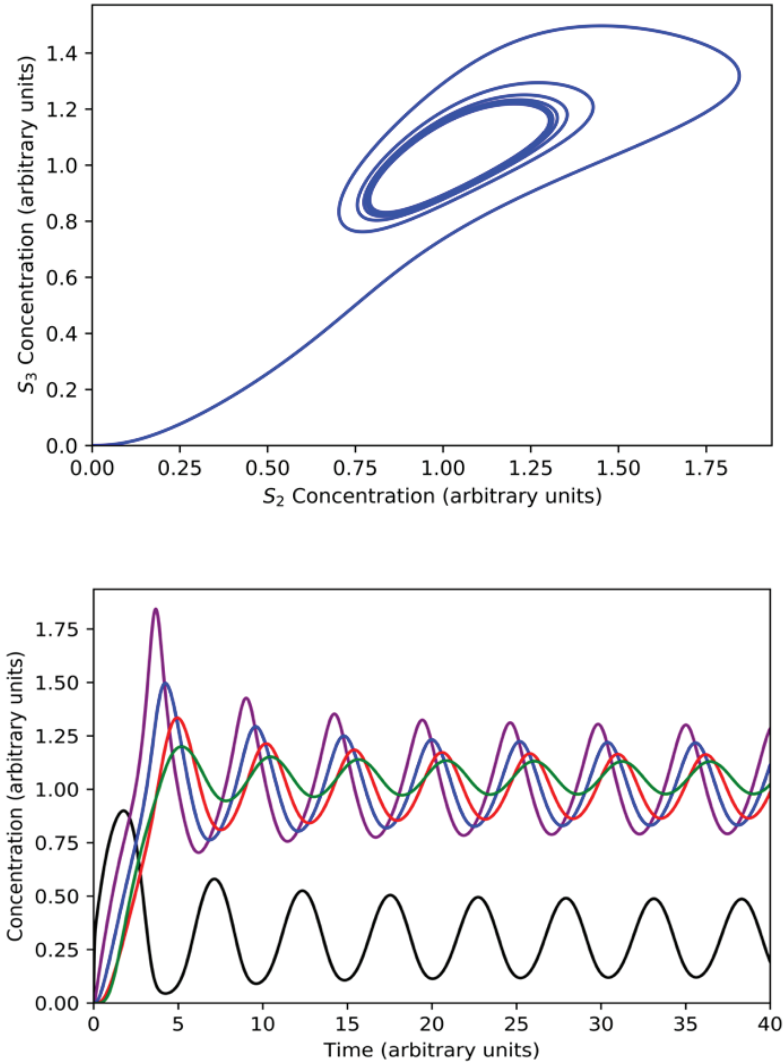
**Note:** The concentration and oscillation amplitude of  $S_4$  are shown to increase by making the last step reversible.

**Table 5.** Oscillation intervals and amplitudes for sustained oscillations with a reversible last step

Metabolites	Oscillation Amplitude	Oscillation Interval
$S_1$	0.33287	0.09472-0.42759
$S_2$	0.79978	0.74726-1.54704
$S_3$	0.26731	0.88438-1.15169
$S_4$	0.19747	2.30410-2.50157
$S_5$	0.03435	1.03506-1.06942

**Note:** It shows an increase in the concentration and amplitude of  $S_4$ .

The model could be easily modified to simulate different metabolic pathways such as branched, cyclic routes, or to implement additional feedback loops (PFBLs or NFBLs) for educational or research purposes. For instance, Figure 13A shows the undamped oscillations observed when  $S_3$ , instead of  $S_2$ , is the activator of  $E_1$  for a  $[S_0]$  equal to 10 a.u. Figure 13B shows the expected limit cycle for  $S_3$  vs.  $S_2$ .



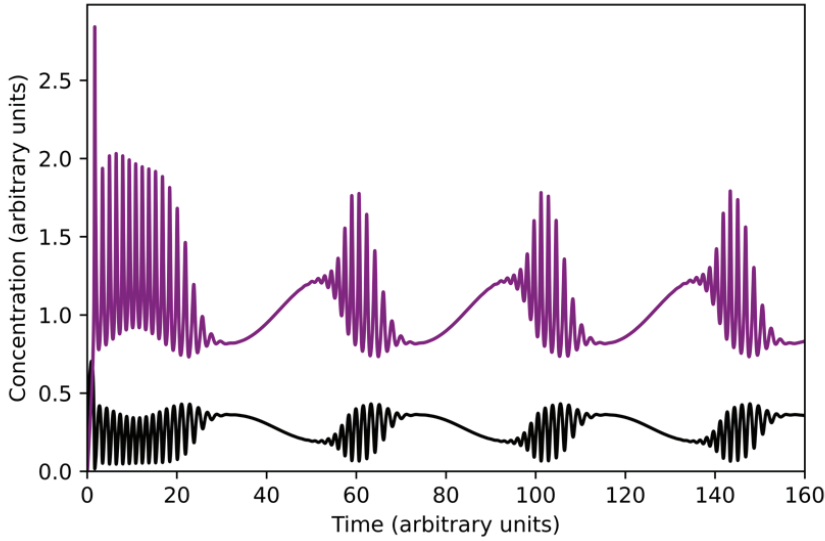
**Figure 13.** (A) Sustained oscillations of the metabolites ( $S_1$ ,  $S_2$ ,  $S_3$ ,  $S_4$ , and  $S_5$ ) when  $S_3$  is the activator of  $E_1$ . (B) Phase plane plot for  $S_3$  vs.  $S_2$  showing a limit cycle.

**Note:** In this experiment, the binomial used to create the PFBL is  $1+S_3^5$  instead of  $1+S_2^3$ .

The effect of changing the constant input flux for a sinusoidal type input flux was also investigated. As previously shown, for constant inlet fluxes, steady states or sustained oscillations are observed. In contrast



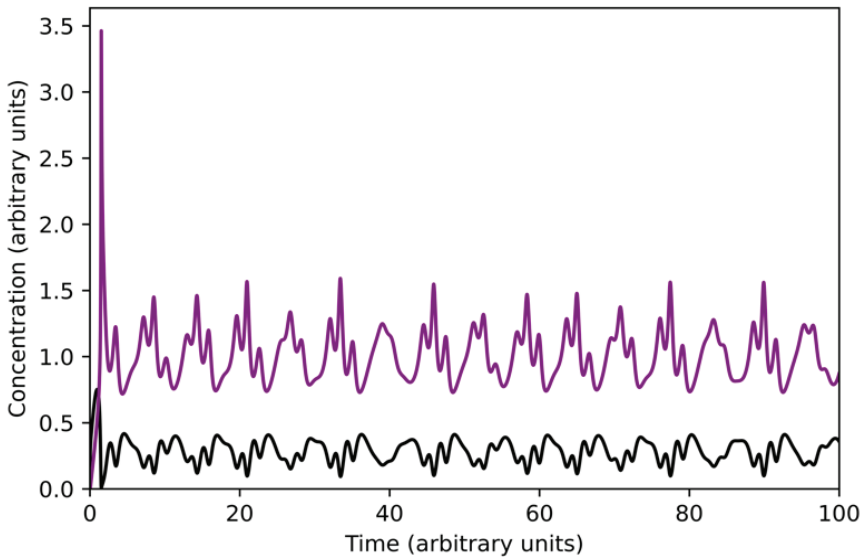
to this, an oscillatory input given by  $S_0 = C + A \sin(\omega t)$  leads to more complex behavior as demonstrated in Figure 14; the value 11.0 a.u. for  $[S_0]$  triggers the sustained oscillations, returning to a non-oscillatory state when  $S_0=9.0$  a.u. which being a periodically excitable system.



**Figure 14.** Complex oscillatory behavior for  $S_1$  and  $S_2$  when a periodic inlet flux, given by  $S_0 = 10 + \sin(0.15t)$ , is used

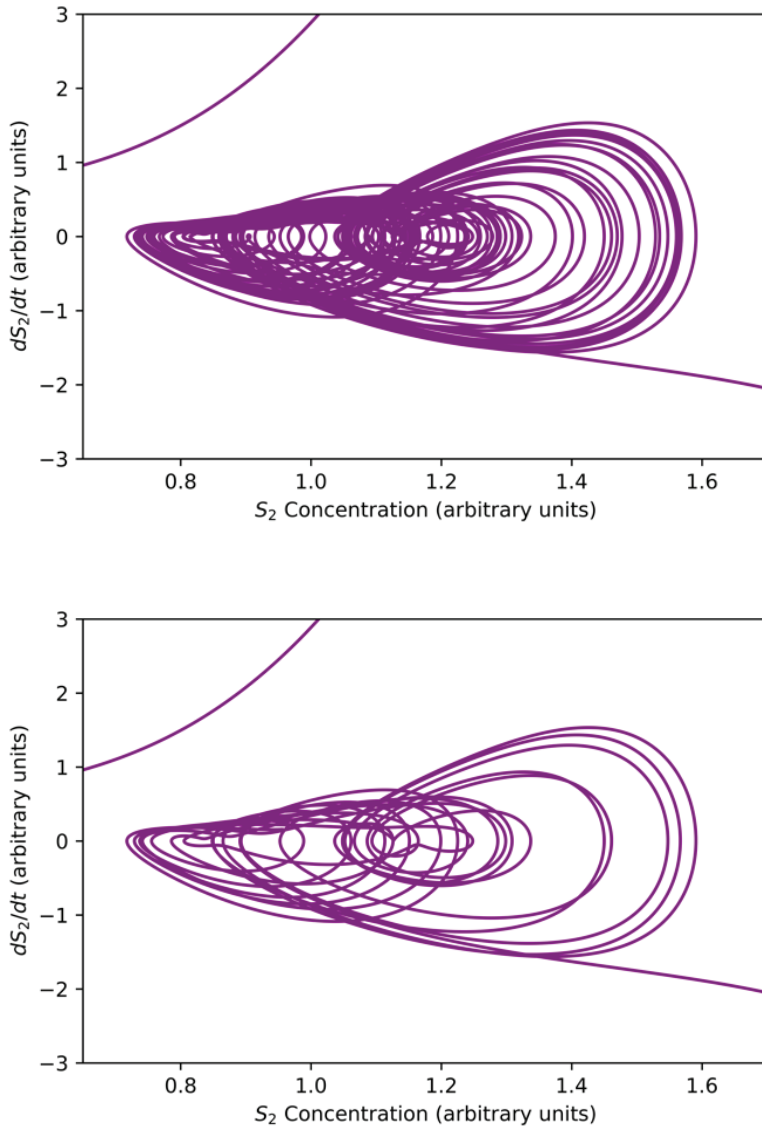
Quasi-chaotic oscillations are observed (Figure 15) when  $C=10$ ,  $A=1$ , and  $\omega=1$ . Figure 16 shows the projections of  $dS_2/dt$  vs.  $S_2$ . Figure 17 shows even more complex and irregular oscillations when an amplitude modulated signal (AM) is used for the inlet flux. Figure 18 shows the phase plane of  $dS_2/dt$  vs.  $S_2$ , which seems to indicate a chaotic dynamic (Suzuki et al., 2016). However, quasiperiodic behavior, or even long period oscillations can result in an ergodically filled torus with a similar appearance to a chaotic attractor (Derrick, 1993). Therefore, although the emergence of chaos as a consequence of a forced oscillatory input in a nonlinear system is well known (Nikolaev et al., 2018), other tests are necessary to ensure its existence, such as the determination of Lyapunov exponents, in addition to the detailed study of the trajectories

in the phase plane. These last two experiments show interesting results (Figures 15 and 17) for developing further studies about the nature of the observed dynamics, derived from the change of the model's original conditions.



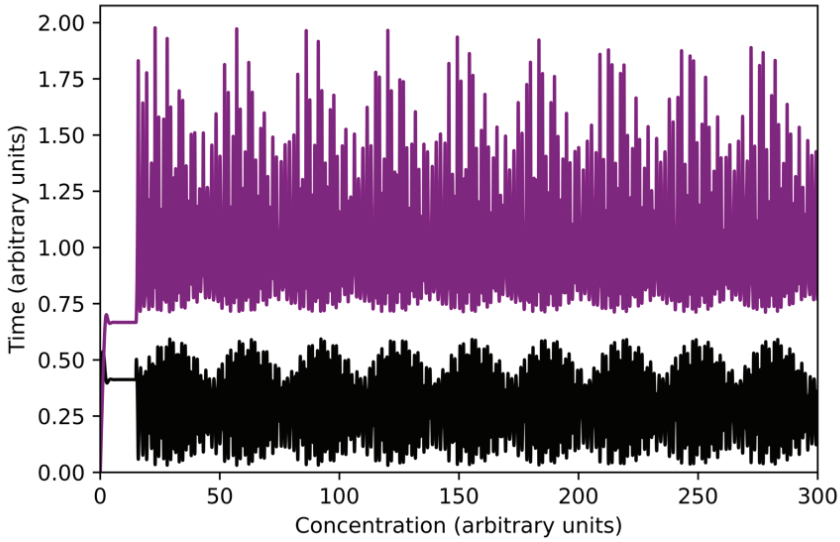
**Figure 15.** Quasi-chaotic oscillations of  $S_1$  and  $S_2$  are observed under a sinusoidal inlet flux

**Note:** For time values higher than 10, the inlet flux is no longer constant but given by the expression  $S_0 = C + A \sin(\omega t)$ .  $C=10$ ,  $A=1$  and  $\omega=1$ . Oscillations in metabolite concentrations return to look similar from time to time, but never identical.



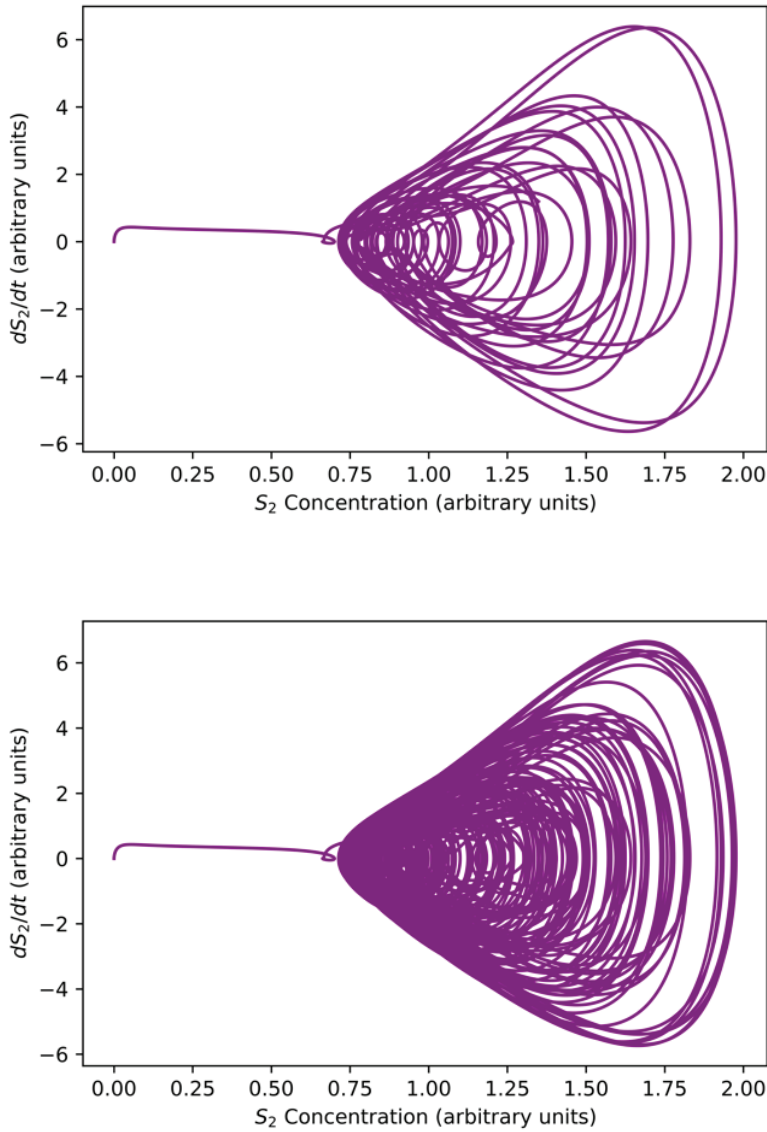
**Figure 16.** Projections of  $dS_2/dt$  vs.  $S_2$  for oscillatory input flux ( $S_0 = C + A \sin(\omega t)$ )

**Note:**  $C=10$ ,  $A=1$  and  $\omega=1$ ).  $t=20$  a.u. (A) and  $t=150$  a.u. (B).



**Figure 17.** Complex oscillatory behavior of  $S_1$  and  $S_2$  observed when an amplitude modulated signal (AM) is used for the inlet flux

**Note:** The AM signal is given by the expression  $S_0 = C (1 + m \cos(\omega_m t) \cos(\omega_p t))$ .  $C=10$ ,  $m=0.45$ ,  $\omega_m=0.10$ ,  $\omega_p=8.75$ . For time values  $< 15$  a.u.,  $[S_0] = 8$  a.u. and a steady state is achieved. When  $t > 15$  u.a., the inlet flux of  $S_0$  becomes oscillatory by the expression given above.



**Figure 18.** Projection of  $dS_2/dt$  vs.  $S_2$  for  $t= 50$  a.u. (A) and  $t= 150$  a.u. (B)

**Note:** According to Suzuki et al., 2016, an irregular figure covering the phase plane is observed instead of the torus expected for a quasiperiodic oscillation, suggesting a chaotic behavior. However, additional tests are required to ensure the existence of chaotic behavior.

#### 4. Conclusions

The behavior of the model was evaluated for  $S_0$  concentrations in the range 0-12 a.u., observing-undamped oscillations for a narrow range of inlet flow values:  $[11.60 > [S_0] > 9.8]$ . Outside this oscillatory interval, the system evolves towards a steady state. The transition between steady states and permanent oscillations take place through damped oscillations. Whilst in the absence of the positive feedback loop the system does not exhibit oscillatory behavior, this is not the only requirement to generate the oscillations since they were shown to depend on enzyme concentrations, in addition to the inlet flow value.

Under a dynamic of sustained oscillations, it was found that the amplitude of the oscillations decreases with the distance to the oscillations source and depends on the value of the kinetic constants and the inlet flux. On the contrary, the phase shift of the oscillations increases with the distance to the oscillophore. Undamped oscillations are also observed when the activator is located further ahead in the sequence of the transformation, and not only when it is the direct product of the core of the oscillatory mechanism. When the system's substrate input is changed from constant to oscillatory, more irregular oscillations emerge that could correspond to chaotic and quasi-chaotic behaviors.

#### Acknowledgments

Biotechnology bachelor students Ceballos M., Díaz K., Hernández D., Hidalgo D., and Martínez P. thank Carolina Juncá for the help in the understanding of Phyton Programming language. The order of authorship obeys only the alphabetical order of the surnames, as each author contributed to the same extent on the realization of this publication.

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